

On partial indices for a matrix Riemann-Hilbert problem

By C. Cercignani and C. E. Siewert, Istituto Matematico del Politecnico di Milano, Milano, Italy

I. Introduction

In a recent paper [1] Siewert and Kelley used a solution, similar to one obtained by Cercignani [2], of the Riemann-Hilbert problem defined by

$$\Phi^+(\mu) = G(\mu) \Phi^-(\mu), \quad \mu \in [0, \infty), \tag{1}$$

where

$$G(\mu) = A^+(\mu)[A^-(\mu)]^{-1}, \tag{2}$$

$$A(z) = I + z \int_{-\infty}^{\infty} \Psi(\mu) \frac{d\mu}{\mu - z}, \tag{3}$$

$$\Psi(\mu) = (\pi)^{-1/2} Q^T(\mu) Q(\mu) e^{-\mu^2} \tag{4}$$

and

$$Q(\mu) = \begin{vmatrix} \left(\frac{2}{3}\right)^{1/2} \left(\mu^2 - \frac{1}{2}\right) & 1 \\ \left(\frac{2}{3}\right)^{1/2} & 0 \end{vmatrix}, \tag{5}$$

to construct a *canonical* solution of the problem. In the process of developing their canonical solution, Siewert and Kelley [1] used the fact, proved previously by Kriese, Chang and Siewert [3], that the partial indices relevant to the considered Riemann-Hilbert problem are $\kappa_1 = \kappa_2 = 1$. Here we wish to show that the partial indices κ_1 and κ_2 can be determined directly from the solutions given by Cercignani [2] and Siewert and Kelley [1]. In order to condense our presentation, we assume that the paper by Siewert and Kelley, to which we hereafter refer as I, is available, and we use, with the inclusion of I in the equation numbers, the results developed there.

II. Analysis

The solution, before it was converted to a *canonical* solution, developed in I is written as

$$\Phi(z) = S^{-1}(z) \begin{vmatrix} U_1(z) & 0 \\ 0 & U_2(z) \end{vmatrix} S(z), \tag{6}$$

where $S(z)$ is given by Eq. (I-12) and $U_\alpha(z)$, $\alpha = 1$ and 2 , are given by Eqs. (I-39) and (I-40). If we let $\Phi_0(z)$ denote a canonical solution then, from I,

$$\Phi_0(z) = [z(z - x_1)]^{-2} \Phi(z) P(z), \tag{7}$$

where $P(z)$ is a matrix of polynomials with $\det P(z) \propto z^2(z - x_1)^2$. From Eq. (7) and I it is apparent that $P(z)$ must be determined so that

$$\Phi(\xi) P(\xi) = 0, \quad \xi = 0 \quad \text{and} \quad \xi = x_1, \tag{8a}$$

and

$$\frac{d}{d\xi} [\Phi(\xi) P(\xi)] = 0, \quad \xi = 0 \quad \text{and} \quad \xi = x_1. \tag{8b}$$

We note from Muskhelishvili [4] that as $|z| \rightarrow \infty$

$$\Phi_0(z) \rightarrow K \text{diag} \{z^{-\kappa_1}, z^{-\kappa_2}\}, \quad \det K \neq 0, \tag{9}$$

where κ_1 and $\kappa_2 \cong \kappa_1$ are the partial indices. We note also that $\kappa_1 + \kappa_2 = \kappa$, where κ can be computed from $\det G(\mu)$; it is also apparent from Eq. (7) that $\kappa = 2$, which agrees with the calculation of Kriese, Chang and Siewert [3].

After using Eqs. (6) and (9) along with the explicit expressions for $S(z)$ and $U_\alpha(z)$ given in I, we can readily deduce from Eq. (7) that as $|z| \rightarrow \infty$

$$P(z) \rightarrow z^3 \begin{vmatrix} a_{11} z^{-\kappa_1} & a_{12} z^{-\kappa_2} \\ a_{21} z^{-\kappa_1} & a_{22} z^{-\kappa_2} \end{vmatrix}. \tag{10}$$

It is clear from Eq. (10) that $\kappa_1 > -2$ for otherwise the second column of $P(z)$ would be identically zero, which is not allowed. From I, we observe that $U_2(z)$ has a double zero at $z = 0$, and thus we conclude from Eqs. (8), after using Eq. (6), that

$$\begin{vmatrix} \left(\frac{3}{2}\right)^{1/2} & 1 \\ 0 & 0 \end{vmatrix} P(0) = 0 \tag{11a}$$

and

$$\begin{vmatrix} \left(\frac{3}{2}\right)^{1/2} & 1 \\ 0 & 0 \end{vmatrix} \frac{d}{d\xi} P(\xi) = 0, \quad \xi = 0. \tag{11b}$$

Now since $U_1(x_1) = 0$, we can conclude from Eq. (8a) that

$$\begin{vmatrix} 0 & 0 \\ 1 & \alpha \end{vmatrix} P(x_1) = 0, \tag{12}$$

where, from I,

$$\alpha = - \left(\frac{1}{6}\right)^{1/2} \left[R(x_1) + x_1^2 + \frac{1}{2} \right]. \tag{13}$$

If we assume for the moment that the elements of the second column of $P(z)$ are linear in z , then Eqs. (11) fix the ratio $P_{12}(z)/P_{22}(z)$ to be $-(2/3)^{1/2}$. From Eq. (12) we thus find a contradiction since $\alpha \neq (2/3)^{1/2}$. It follows that the second column of $P(z)$ must be at least quadratic in z and therefore that $\kappa_1 = \kappa_2 = 1$, which agrees with the conclusion of Kriese, Chang and Siewert [3].

III. Concluding remarks

It is worthwhile to recall from the work of Muskhelishvili [4] that the partial indices play a crucial role in the theory of systems of singular-integral equations. Muskhelishvili [4] in fact uses these indices to determine if a given system can be solved, and if so the number and type of constraints that may be required to yield a unique solution. Though there are problems in regard to coupled systems in particle transport theory, see for example [3, 5, 6], for which the partial indices have been determined, there are also important cases, even for two-vector problems, where these indices remain elusive [7, 8]. In our continuing work we intend to see if the idea developed here can be used to determine the partial indices for Riemann-Hilbert problems more difficult than the one considered at present.

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References

- [1] C. E. Siewert and C. T. Kelley, *Z. angew. Math. Phys.* 31, 344 (1980).
- [2] C. Cercignani, *Trans. Theory and Stat. Phys.* 6, 29 (1977).
- [3] J. T. Kriese, T. S. Chang, and C. E. Siewert, *Int. J. Eng. Sci.* 12, 441 (1974).
- [4] N. I. Muskhelishvili, *Singular integral equations*. Noordhoff, Groningen, Holland (1953).
- [5] C. E. Siewert and E. E. Burniston, *The Astrophys. J.* 174, 629 (1972).
- [6] C. E. Siewert, E. E. Burniston, and J. T. Kriese, *J. Nucl. Energy* 26, 469 (1972).
- [7] J. R. Thomas, Jr. and C. E. Siewert, *Trans. Theory and Stat. Phys.* 8, 219 (1979).
- [8] C. E. Siewert and J. R. Thomas, Jr., *Z. angew. Math. Phys.* 33, 202–218 (1982).

Abstract

An analytical solution of the matrix Riemann-Hilbert problem relevant to the BGK model in the field of rarefied-gas dynamics is used to deduce the partial indices basic to a canonical solution of the considered Riemann-Hilbert problem.

Sommario

Si fa uso di una soluzione analitica del problema di Riemann-Hilbert che origina dal modello BGK nel campo della dinamica dei gas rarefatti per dedurre gli indici necessari per ottenere una soluzione canonica del problema di Riemann-Hilbert.

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