

NOTE

A NOTE ON RADIATIVE TRANSFER

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Abstract—The formulation of radiative transfer problems with internal energy sources is discussed

1 DISCUSSION

In a recent publication Yuen and Wong¹ suggested that one of our early works² has only limited applicability in heat-transfer calculations because energy conservation was not considered and because a temperature distribution in the medium was assumed. Here we would like to point out that, with the exception of the limitation to isotropic scattering and from a mathematical point-of-view, the formulation we considered in that work² is sufficiently general so as to include the Yuen-Wong problem as a special case. To be general, we start with a formulation that includes *L*th order anisotropic scattering and that allows both specular and diffuse reflection at the surfaces. Consider

$$\mu \frac{\partial}{\partial \tau} I(\tau, \mu) + I(\tau, \mu) = \frac{\omega}{2} \sum_{l=0}^L \beta_l P_l(\mu) \int_{-1}^1 P_l(\mu') I(\tau, \mu') d\mu' + S(\tau) \quad (1)$$

and the boundary conditions

$$I(L, \mu) = F_1(\mu) + \rho_1^s I(L, -\mu) + 2\rho_1^d \int_{-1}^1 I(L, -\mu') \mu' d\mu', \quad \mu > 0, \quad (2a)$$

and

$$I(R, -\mu) = F_2(\mu) + \rho_2^s I(R, \mu) + 2\rho_2^d \int_{-1}^1 I(R, \mu') \mu' d\mu', \quad \mu > 0, \quad (2b)$$

where $F_1(\mu)$ and $F_2(\mu)$ are given. This problem, as recently reported,³ can be solved for a given $S(\tau)$.

We now investigate the type of problem considered by Yuen and Wong,¹ i.e.

$$\mu \frac{\partial}{\partial \tau} \psi(\tau, \mu) + \psi(\tau, \mu) = \frac{\omega}{2} \sum_{l=0}^L \gamma_l P_l(\mu) \int_{-1}^1 P_l(\mu') \psi(\tau, \mu') d\mu' + (1 - \omega) I_b(\tau), \quad (3)$$

$$\psi(L, \mu) = F_1(\mu) + \rho_1^s \psi(L, -\mu) + 2\rho_1^d \int_0^1 \psi(L, -\mu') \mu' d\mu', \quad \mu > 0, \quad (4a)$$

$$\psi(R, -\mu) = F_2(\mu) + \rho_2^s \psi(R, \mu) + 2\rho_2^d \int_0^1 \psi(R, \mu') \mu' d\mu', \quad \mu > 0, \quad (4b)$$

and

$$\frac{d}{d\tau} \int_{-1}^1 \psi(\tau, \mu') \mu' d\mu' = K(\tau), \quad (5)$$

where $F_1(\mu)$, $F_2(\mu)$ and the internal generation $K(\tau)$ are assumed given, but $I_b(\tau)$ is to be determined. Equation (5) represents the requirement that the energy generated in the medium be dissipated only by the radiation process. If we integrate Eq (3) over μ from -1 to 1 we find

$$\frac{d}{d\tau} \int_{-1}^1 \psi(\tau, \mu) \mu d\mu + (1 - \omega) \int_{-1}^1 \psi(\tau, \mu) d\mu = 2(1 - \omega)I_b(\tau) \quad (6)$$

or, after we use Eq (5),

$$2(1 - \omega)I_b(\tau) = K(\tau) + (1 - \omega) \int_{-1}^1 \psi(\tau, \mu) d\mu \quad (7)$$

With Eq (7) we can eliminate the unknown $I_b(\tau)$ from Eq (3) to obtain

$$\mu \frac{\partial}{\partial \tau} \psi(\tau, \mu) + \psi(\tau, \mu) = \frac{1}{2} \int_{-1}^1 \psi(\tau, \mu') d\mu' + \frac{1}{2} \sum_{l=1}^L \omega \gamma_l P_l(\mu) \int_{-1}^1 P_l(\mu') \psi(\tau, \mu') d\mu' + \frac{1}{2} K(\tau), \quad (8)$$

which is to be solved subject to the boundary conditions given by Eqs (4)

It is therefore clear that in Eq (1) we can let $2S(\tau) = K(\tau)$, use $\beta_l = \omega \gamma_l$, $l \geq 1$, and set $\omega = 1$ to obtain a problem of the type considered by Yuen and Wong;¹ thus our formulation as given by Eqs (1) and (2) has enhanced applicability rather than limited utility in heat-transfer calculations

REFERENCES

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- 3 C E Siewert, J R Maiorino, and M N Ozisik, *JQSRT* **23**, 565 (1980)