

A Computation of the X and Y Functions for a Nongrey Model with Complete Frequency Redistribution

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In this presentation we discuss a computation of the classic X and Y functions for use in the solution of a class of nongrey radiative-transfer problems defined for finite layers. We consider a model that allows for scattering with complete frequency redistribution (completely noncoherent scattering) and continuum absorption. In order to illustrate basic numerical aspects of our solution for the X and Y functions, we consider some test problems based on the Doppler and the Lorentz forms of the line-scattering coefficient.

Our problem is based on the equation of transfer written as

$$\mu \frac{\partial}{\partial \tau} I_{\pm}(\tau, \mu) = [\phi(x) + \beta][S_{\pm}(\tau) - I_{\pm}(\tau, \mu)]$$

where $S_{\pm}(\tau)$ is the source function,

$$[\phi(x) + \beta]S_{\pm}(\tau) = \frac{1}{2} \varpi \phi(x) \int_{-\infty}^{\infty} \phi(x') \int_{-1}^1 I_{\pm}(\tau, \mu') d\mu' dx' + [\rho\beta + (1 - \varpi)\phi(x)]B(\tau),$$

and $B(\tau)$ is the Planck function. Here $\tau \in [0, \tau_0]$ is the optical variable, τ_0 is the optical thickness of the considered layer and $\mu \in [-1, 1]$ is the cosine of the polar angle (as measured from the positive τ axis) that describes the direction of propagation of the radiation. In addition $\varpi \in [0, 1]$ is the albedo for single scattering, $\beta \geq 0$ is the ratio of the continuum absorption coefficient to the average line coefficient, ρ is the ratio of the continuum source function to the Planck function and $\phi(x)$ is the line-scattering profile. We note that radiative-transfer problems in finite plane-parallel media are often expressed in terms of X and Y functions, and so our goal here is to compute those functions for this nongrey radiative-transfer model.

We consider the X and Y functions to be defined by the coupled singular-integral equations

$$\lambda(\eta)X(\eta) - \eta \int_0^{\gamma} \Psi(\xi)X(\xi) \frac{d\xi}{\xi - \eta} + \eta e^{-\tau_0/\eta} \int_0^{\gamma} \Psi(\xi)Y(\xi) \frac{d\xi}{\xi + \eta} = 1$$

and

$$\lambda(\eta)Y(\eta) - \eta \int_0^{\gamma} \Psi(\xi)Y(\xi) \frac{d\xi}{\xi - \eta} + \eta e^{-\tau_0/\eta} \int_0^{\gamma} \Psi(\xi)X(\xi) \frac{d\xi}{\xi + \eta} = e^{-\tau_0/\eta}$$

for $\eta \in [0, \gamma]$. We note that we use the symbol \int to denote that integrals are to be evaluated in the Cauchy principal-value sense. In addition $\gamma = 1/\beta$, $\Psi(\xi)$ is the characteristic function and

$$\lambda(\eta) = 1 + \eta \int_{-\infty}^{\gamma} \Psi(\xi) \frac{d\xi}{\xi - \eta}.$$

In this work we choose to use the F_N method to compute numerical solutions to Eqs. (3). We thus expand the X and Y functions in terms of a set of basis functions, substitute these approximations into Eqs. (3), consider the resulting equations at selected values of $\eta \in [0, \gamma]$ and solve a system of linear algebraic equations to find the required expansion coefficients in the approximate solutions. We note that the challenges of computing the desired X and Y functions for the considered radiative-transfer model are not insignificant!