

ON PARTICULAR SOLUTIONS OF THE EQUATION OF TRANSFER

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Abstract—The particular solutions corresponding to two very special inhomogeneous terms of the equation of transfer with linearly anisotropic scattering are reported

I INTRODUCTION

It is clear that the analysis of radiative transfer problems, based on the inhomogeneous equation

$$\mu \frac{\partial}{\partial \tau} I(\tau, \mu) + I(\tau, \mu) = \frac{\omega}{2} \int_{-1}^1 (1 + b\mu\mu') I(\tau, \mu') d\mu' + S(\tau), \quad (1)$$

can proceed along classical lines⁽¹⁾ once a particular solution $I_p(\tau, \mu)$, corresponding to a given inhomogeneous source term $S(\tau)$, is established. In this paper, we wish to report two particular solutions which correspond to source terms explicitly omitted in an existing tabulation⁽²⁾ of particular solutions, namely

$$S(\tau) = e^{-\tau/\eta_0} \quad \text{and} \quad S(\tau) = \tau e^{-\tau/\eta_0} \quad (2)$$

where η_0 denotes a “discrete eigenvalue” of the homogeneous equation of transfer:

$$\Lambda(\eta_0) = \omega R(\eta_0, \eta_0) \left[1 - \eta_0 \tanh^{-1} \frac{1}{\eta_0} \right] + 1 - \omega = 0, \quad (3)$$

where

$$R(x, y) = 1 + b(1 - \omega)xy \quad (4)$$

II ANALYSIS

If we first consider $S(\tau) = \exp(-\tau/\eta_0)$ and substitute

$$I_p(\tau, \mu) = \varphi(\eta_0, \mu) e^{-\tau/\eta_0} [A\tau + B(\mu)], \quad (5)$$

where

$$\varphi(\eta_0, \mu) = \frac{\omega}{2} \eta_0 \frac{R(\eta_0, \mu)}{\eta_0 - \mu}, \quad (6)$$

into eqn (1), we find that the constant A and the unknown $B(\mu)$ should satisfy

$$A\mu\varphi(\eta_0, \mu) + \frac{1}{2} \omega R(\eta_0, \mu) B(\mu) = 1 + \frac{\omega}{2} \int_{-1}^1 B(\mu') \varphi(\eta_0, \mu') (1 + b\mu\mu') d\mu'. \quad (7)$$

Multiplying eqn (7) by $\varphi(\eta_0, \mu)$ and integrating over μ from -1 to 1 , we find immediately that

$$\frac{1}{A} = N(\eta_0) = \int_{-1}^1 \mu \varphi^2(\eta_0, \mu) d\mu. \quad (8)$$

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To find $B(\mu)$, we can first integrate eqn (7) from -1 to 1 to obtain

$$A\eta_0(1-\omega) + \frac{1}{2}\omega \int_{-1}^1 R(\eta_0, \mu)B(\mu) d\mu = 2 + \omega B_0, \quad (9)$$

where, in general,

$$B_\alpha = \int_{-1}^1 \varphi(\eta_0, \mu)B(\mu)\mu^\alpha d\mu. \quad (10)$$

Thus, after noting that

$$B_1 = \eta_0 B_0 - \frac{1}{2}\omega\eta_0 \int_{-1}^1 R(\eta_0, \mu)B(\mu) d\mu, \quad (11)$$

we can write eqn (7) as

$$A\mu\varphi(\eta_0, \mu) + \frac{1}{2}\omega R(\eta_0, \mu)B(\mu) = 1 + \frac{1}{2}\omega R(\eta_0, \mu)B_0 + \frac{1}{2}\omega b\eta_0\mu[-2 + A\eta_0(1-\omega)]. \quad (12)$$

Since $\varphi(\eta_0, \mu) \exp(-\tau/\eta_0)$ is a solution of the homogeneous version of eqn (1), it is clear that the constant B_0 appearing in eqn (12) can be chosen arbitrarily. We now take $B_0 = -2/\omega$ to obtain the final result

$$B(\mu) = \frac{-2}{\omega R(\eta_0, \mu)} \left\{ A\mu\varphi(\eta_0, \mu) + b\eta_0\mu \left[1 - \frac{1}{2}\omega A\eta_0(1-\omega) \right] \right\}. \quad (13)$$

The algebra associated with the case $S(\tau) = \tau \exp(-\tau/\eta_0)$ is considerably more tedious, and thus we will only list here the final results:

$$I_p(\tau, \mu) = \varphi(\eta_0, \mu) e^{-\tau/\eta_0} [C\tau^2 + D(\mu)\tau + E(\mu)], \quad (14)$$

where

$$C = \frac{1}{2N(\eta_0)}, \quad (15)$$

$$D(\mu) = -\frac{4C}{\omega R(\eta_0, \mu)} \left[\mu\varphi(\eta_0, \mu) - \omega R(\eta_0, \mu) \left(\frac{L(\eta_0)}{N(\eta_0)} - K(\eta_0) \right) - N(\eta_0) \right. \\ \left. \times \{1 - \omega b\eta_0\mu[1 - C\eta_0(1-\omega)]\} \right], \quad (16)$$

and

$$E(\mu) = \frac{2}{\omega R(\eta_0, \mu)} \left[-\mu\varphi(\eta_0, \mu)D(\mu) + \frac{\omega}{2} b\eta_0\mu R \right]. \quad (17)$$

Here we have let

$$L(\eta_0) = \int_{-1}^1 \frac{\varphi^3(\eta_0, \mu)}{\omega R(\eta_0, \mu)} \mu^2 d\mu, \quad (18)$$

$$K(\eta_0) = \int_{-1}^1 \frac{\varphi^2(\eta_0, \mu)}{\omega R(\eta_0, \mu)} \{1 + \omega b\eta_0\mu[-1 + C\eta_0(1-\omega)]\} \mu d\mu, \quad (19)$$

and

$$R = \frac{\eta_0(1-\omega)}{N(\eta_0)} \left[\eta_0 - \left(\frac{2}{1-\omega} \right) N(\eta_0) + 2 \frac{L(\eta_0)}{N(\eta_0)} - 2K(\eta_0) \right]. \quad (20)$$

We note that Ref. (2) contains particular solutions for inhomogeneous source terms of the form $\exp(-\tau/\eta)$ and $\tau \exp(-\tau/\eta)$, $\eta \neq \eta_0$, and thus the results given here complete the tabulation of Ref. (2).

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REFERENCES

- 1 K M CASE and P F ZWEIFEL, *Linear Transport Theory* Addison-Wesley, Reading, Mass (1967)
- 2 M N OZISIK and C E SIEWERT, *Nucl Sci Engng* **40**, 491 (1970)